

波浪中の増加抵抗公式について(末)

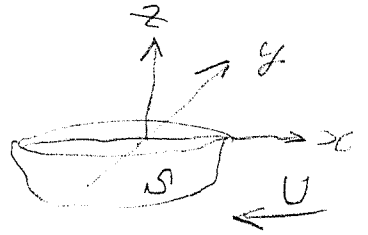
概要

平面波のエネルギー、運動量の式から漸近展開を用いて波浪中の増加抵抗を初等的に説明する。

1. 漸近展開

水面条件は

$$\left. \frac{\partial \phi}{\partial z} \right|_{z=0} = -(i\omega - U \frac{\partial}{\partial x}) \zeta(x, y), \quad (1.1)$$



$$\zeta(x, y) = \frac{1}{g} (i\omega - U \frac{\partial}{\partial x}) \phi \Big|_{z=0}, \quad (1.2)$$

$$(i\omega - U \frac{\partial}{\partial x})^2 \phi + g \frac{\partial \phi}{\partial z} = 0, \quad z=0, \quad (1.3)$$

核函数を

$$S(x, y, z) = \frac{1}{4\pi} \left(\frac{1}{r_1} - \frac{1}{r_2} \right) - \frac{g}{4\pi^2} \int_{-\pi}^{\pi} \int_0^{\infty} \frac{e^{-kz} e^{ik[(x-x')\cos\theta + (y-y')\sin\theta]}}{A(k, \theta)} k dk d\theta, \quad (1.4)$$

$$A(k, \theta) = (\omega + kU\cos\theta)^2 - gk - i\mu(kU\cos\theta + \omega), \quad (1.5)$$

とおくと

$$\phi(p) = \iint_{S+\bar{H}} \left\{ \frac{\partial \phi}{\partial n} S - \phi \frac{\partial S}{\partial n} \right\} dS, \quad (1.6)$$

のより核函数を表現が出来るから、今

$$H(k, \theta) = \frac{1}{4\pi} \iint_{S+\bar{H}} \left(\frac{\partial \phi}{\partial n} - \phi \frac{\partial}{\partial n} \right) e^{ik[(x-x')\cos\theta + (y-y')\sin\theta] + kz} dS, \quad (1.7)$$

とおくと

$$\phi(p) = \phi_0(p) + \phi_1(p), \quad (1.8)$$

$$\phi_0(p) = \frac{1}{4\pi} \iint_{S+\bar{H}} \left(\frac{\partial \phi}{\partial n} - \phi \frac{\partial}{\partial n} \right) \left(\frac{1}{r_1} - \frac{1}{r_2} \right) dS, \quad (1.9)$$

$$\phi_1(\rho) = -\frac{g}{\pi} \int_{-\pi}^{\pi} \int_0^{\infty} \frac{H(k, \theta) e^{-kz} e^{-i(kx \cos \theta + y \sin \theta)}}{A(k, \theta)} k dk d\theta, \quad (1.10)$$

と書ける。

漸近近傍を考慮する時は ϕ_1 のみとすればよい。

またこの時 (1.2) により

$$\zeta(x, y) \doteq \frac{1}{\pi i} \int_{-\pi}^{\pi} \int_0^{\infty} \frac{H(k, \theta) e^{-kz} e^{-i(kx \cos \theta + y \sin \theta)}}{A(k, \theta)} k dk d\theta, \quad (1.11)$$

さて $A(k, \theta) = U^2 \cos^2 \theta (k - k_1)(k - k_2), \quad (1.12)$

$$\begin{aligned} \zeta(x, y) \doteq & \frac{2}{U^2} \int_{\frac{\pi}{2}}^{\frac{3\pi}{2}} \frac{H(k_1, \theta) e^{-i k_1 y \cos(\theta - \varphi)}}{(k_1 - k_2) \cos^2 \theta} d\theta \\ & + \frac{2}{U^2} \int_{\frac{\pi}{2}}^{\frac{3\pi}{2}} \frac{H(k_2, \theta) e^{-i k_2 y \cos(\theta - \varphi)}}{(k_2 - k_1) \cos^2 \theta} d\theta, \quad (1.13) \end{aligned}$$

$$U^2 \cos^2 \theta (k_2 - k_1) = g \sqrt{1 - \frac{4\omega^2 U^2 \cos^2 \theta}{g^2}},$$

定常点とは

$$f(\theta) = k \cos(\theta - \varphi) \quad (1.14)$$

とすると $f'(\theta) = 1 [k' \cos(\theta - \varphi) - k \sin(\theta - \varphi)] = 0 \quad (1.15)$

$$\therefore \tan(\theta - \varphi) = \frac{k'(\theta)}{k(\theta)} \quad (1.16)$$

次に $f''(\theta) = 1 [(k'' - k) \cos(\theta - \varphi) - 2k' \sin(\theta - \varphi)]$,

とすると (1.16) を代入すると

$$f''(\theta) = k \cos(\theta - \varphi) \left[\left(\frac{k'}{k}\right)' - 1 - \frac{k'^2}{k^2} \right] \quad (1.17)$$

と得る。

ここで (1.16) を θ で微分すると (今は φ は θ の関数となる)

$$\begin{aligned} \left(\frac{k'}{k}\right)' &= \sec^2(\theta - \varphi) \left[1 - \frac{d\varphi}{d\theta}\right] \\ &= 1 + \frac{k'^2}{k^2} - \sec^2(\theta - \varphi) \frac{d\varphi}{d\theta} \end{aligned} \quad (1.18)$$

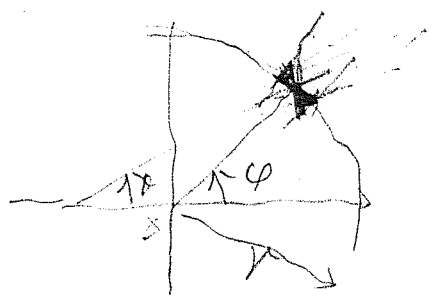
これを (1.17) に代入すると (Newman)

$$f''(\theta) = - \frac{k}{\cos(\theta - \varphi)} \frac{d\varphi}{d\theta} \quad (1.19)$$

したがって

$$\xi = \frac{2\sqrt{2\pi}}{g} \frac{1}{r} \sum_{j,n} \pm \frac{H(k_{jn} \theta_n) k_{jn} \omega_{jn} \rho^{-i k_{jn} \cos(\theta_n - \varphi) \pm \frac{\pi}{2} i}}{\sqrt{|f''(\theta_n)|} \sqrt{1 - \frac{4\omega \nu \cos \theta_n}{g}}} \quad (1.20)$$

無限大の円筒面の面積素 $r d\varphi$ を
 単位時間には
 通じて出てゆく位のエネルギーは



$$dW = \frac{P_0}{2} \left(\frac{c_j}{2} - U \cos \theta \right) \iint \rho \cos(\theta - \varphi) r d\varphi$$

運動量は

$$dM = \frac{P_0}{2} \left(\frac{1}{2} - \frac{U \cos \theta}{c_j} \right) \iint \rho \cos(\theta - \varphi) r d\varphi$$

$$dW = \frac{P_0}{2} \left(\frac{P_0}{2} - \frac{U \cos \theta}{c_j} \right) \frac{8\pi}{g^2} \frac{|H(k; \theta)|^2 k_r^2 \omega_j^2 d\theta}{\left(1 - \frac{4\omega U \cos \theta}{g}\right)}$$

$$= \frac{2\pi P_0 |H|^2 k_r^2 d\theta}{\sqrt{1 - \frac{4\omega U \cos \theta}{g}}}$$

$$dW = 2\pi P_0 \frac{|H|^2 k_j \omega_j d\theta}{\sqrt{1 - \frac{4\omega U \cos \theta}{g}}}$$

$$k_j c_j = \omega_j$$

$$\omega_j = \omega + k_j U \cos \theta$$

x 方向の運動量は $dM \cos \theta$.

→ 部分エネルギー

$$dW = dE + (dR)U$$

← 振子の平均位置

$$dE = 2\pi P_0 \omega \frac{|H|^2 k_r d\theta}{\sqrt{1 - \frac{4\omega U \cos \theta}{g}}}$$

附錄 A

$$\begin{aligned}
 A(k, \theta) &= (\omega + kU \cos \theta)^2 - gk - i\mu(\omega + kU \cos \theta) \\
 &= k^2 U^2 \cos^2 \theta + k(2\omega U \cos \theta - g - i\mu U \cos \theta) + \omega^2 - i\mu \omega,
 \end{aligned} \tag{A.1}$$

$$\left. \begin{matrix} k_1 \\ k_2 \end{matrix} \right\} = \frac{+g - 2\omega U \cos \theta + i\mu U \cos \theta \mp g \sqrt{1 - \frac{4U \cos \theta (\omega - \frac{i\mu}{2})}{g}}}{2U^2 \cos^2 \theta}, \tag{A.2}$$

$$k_j U \cos \theta + \omega = \omega_j, \quad j=1, 2, \tag{A.3}$$

$$\left. \begin{matrix} \omega_1 \\ \omega_2 \end{matrix} \right\} = \frac{g + i\mu U \cos \theta \mp g \sqrt{1 - \frac{4U \cos \theta}{g}}}{2U \cos \theta}, \tag{A.4}$$

$$\omega_j^2 / g = k_j^2 = g / c_j^2, \quad \frac{g}{c_j} = \omega_j, \tag{A.5}$$

$$\left. \begin{matrix} c_1 \\ c_2 \end{matrix} \right\} = \frac{2U \cos \theta}{1 \mp \sqrt{1 - \frac{4U \cos \theta}{g}}}, \tag{A.6}$$

$$\left. \begin{matrix} \frac{c_1}{2} > U \cos \theta \\ \frac{c_2}{2} < U \cos \theta \end{matrix} \right\} \begin{matrix} \text{前進} \\ \text{不前進} \end{matrix} \text{ for } \cos \theta > 0. \tag{A.7}$$

$$\left. \begin{matrix} \frac{c_1}{2} = \frac{+U |\cos \theta|}{\sqrt{1 + \frac{4U |\cos \theta|}{g}} - 1} > 0, \text{ for } \cos \theta < 0, \\ \frac{c_2}{2} = \frac{-U |\cos \theta|}{\sqrt{1 + \frac{4U |\cos \theta|}{g}} + 1} < 0, \left| \frac{c_2}{2} \right| < U |\cos \theta|, \text{ 不前進} \end{matrix} \right\} \text{ for } \cos \theta < 0$$

$$\begin{aligned}
 (k_2 - k_1) U^2 \cos^2 \theta &= g \sqrt{1 - \frac{4U \cos \theta}{g}}, \\
 &= (\omega_2 - \omega_1) U \cos \theta,
 \end{aligned} \tag{A.8}$$

$$\begin{aligned}
 (k_1 + k_2) U^2 \cos^2 \theta &= g - 2\omega U \cos \theta \\
 k_1 k_2 U^2 \cos^2 \theta &= \omega^2 \\
 (\omega_1 + \omega_2) U \cos \theta &= g, \quad \omega_1 \omega_2 = U \cos \theta = g \omega
 \end{aligned}
 \quad \left. \vphantom{\begin{aligned} (k_1 + k_2) U^2 \cos^2 \theta &= g - 2\omega U \cos \theta \\ k_1 k_2 U^2 \cos^2 \theta &= \omega^2 \\ (\omega_1 + \omega_2) U \cos \theta &= g, \quad \omega_1 \omega_2 = U \cos \theta = g \omega \end{aligned}} \right\} (A.9)$$

$$\frac{1}{c_1} + \frac{1}{c_2} = \frac{1}{U \cos \theta}$$

$$\sqrt{1 - \frac{4\omega U \cos \theta}{g}} = \frac{2U \cos \theta}{c_2} - 1 = 1 - \frac{2U \cos \theta}{c_1} \quad (A.10)$$

$$\frac{k'_j}{k_j} = 2 \frac{\omega'_j}{\omega_j}, \quad \left(\frac{k'_j}{k_j} \right)' = 2 \left\{ \frac{\omega''_j}{\omega_j} - \left(\frac{\omega'_j}{\omega_j} \right)^2 \right\} \quad (A.11)$$

$$\begin{aligned}
 \frac{\omega'_j}{\omega_j} &= \frac{1 \mp \frac{2\omega U \cos \theta}{g}}{\left(1 \mp \sqrt{1 - \frac{4\omega U \cos \theta}{g}} \right) \sqrt{1 - \frac{4\omega U \cos \theta}{g}}} + \tan \theta = \left(1 \mp \frac{\omega (c_j + 1)}{g \sqrt{1 - \frac{4\omega U \cos \theta}{g}}} \right) \tan \theta \\
 &= \left(1 \mp \frac{\omega}{\omega_j \sqrt{1 - \frac{4\omega U \cos \theta}{g}}} \right) \tan \theta \\
 &= \left(1 \mp \frac{\omega g}{\omega_j (\omega_2 - \omega_1) U \cos \theta} \right) \tan \theta \\
 &= \frac{1 \mp \tan \theta}{U \omega_j (\omega_2 - \omega_1)} \left\{ \frac{\omega_j (\omega_2 - \omega_1) U \cos \theta}{g} \mp \omega \right\}
 \end{aligned}$$

$$\therefore \frac{\omega'_1}{\omega_1} = \frac{g \tan \theta}{U \omega_1 (\omega_2 - \omega_1)} (\omega - \omega_1) = - \frac{\omega_1}{\omega_2 - \omega_1} \tan \theta \quad \left. \vphantom{\frac{\omega'_1}{\omega_1}} \right\} (A.12)$$

$$\frac{\omega'_2}{\omega_2} = \frac{\omega_2}{\omega_2 - \omega_1} \tan \theta$$

$$\begin{aligned}
 \left(\frac{\omega_j'}{\omega_j}\right)' &= \frac{\omega_j'}{\omega_j - \omega_{j+1}} \sec^2 \theta + \tan \theta \left(\frac{\omega_j'}{\omega_j - \omega_{j+1}} - \frac{\omega_j'(\omega_j' - \omega_{j+1}')}{(\omega_j - \omega_{j+1})^2} \right) \\
 &= \frac{\omega_j' \sec^2 \theta}{\omega_j - \omega_{j+1}} + \frac{\tan \theta \omega_j \omega_{j+1}'}{(\omega_j - \omega_{j+1})^2} \left(\omega_{j+1}' - \frac{\omega_j' \omega_{j+1}'}{\omega_j} \right) \\
 &= \frac{\omega_j' \sec^2 \theta}{\omega_j - \omega_{j+1}} - \frac{\tan^2 \theta (\omega_j + \omega_{j+1}')}{(\omega_j - \omega_{j+1})^3} \omega_j \omega_{j+1}' + \frac{\tan^2 \theta \omega_j \omega_{j+1}'}{(\omega_j - \omega_{j+1})^3} \left[\frac{(\omega_j' - \omega_{j+1}')^2}{\omega_j - \omega_{j+1}'} - \omega_{j+1}' (\omega_j + \omega_{j+1}') \right] \\
 &= \frac{\omega_j'}{\omega_j - \omega_{j+1}} + \frac{\omega_j' \tan^2 \theta}{(\omega_j - \omega_{j+1})^3} (\omega_j^2 - 3\omega_j \omega_{j+1}') \quad , \quad (A.13)
 \end{aligned}$$

$$\begin{aligned}
 \left(\frac{k'}{k}\right)' - \left(1 + \frac{k'^2}{k^2}\right) &= 2 \left(\frac{\omega_j'}{\omega_j}\right)' - \left\{1 + 4 \left(\frac{\omega_j'}{\omega_j}\right)^2\right\} \\
 &= 2 \left[\frac{\omega_j'}{\omega_j - \omega_{j+1}} + \frac{\omega_j' \tan^2 \theta}{(\omega_j - \omega_{j+1})^3} (\omega_j^2 - 3\omega_j \omega_{j+1}') \right] - 1 - \frac{4\omega_j'^2 \tan^2 \theta}{(\omega_j - \omega_{j+1}')^2} \\
 &= \frac{\omega_j + \omega_{j+1}'}{\omega_j - \omega_{j+1}} + \frac{2\omega_j'^2 \tan^2 \theta}{(\omega_j - \omega_{j+1})^3} (\omega_j - 3\omega_{j+1}' - 2(\omega_j - \omega_{j+1}')) \\
 &= \frac{\omega_j + \omega_{j+1}'}{\omega_j - \omega_{j+1}} \left[1 + \frac{2\omega_j'^2 \tan^2 \theta}{(\omega_j - \omega_{j+1}')^2} \right] \quad , \quad (A.14)
 \end{aligned}$$

$$\begin{aligned}
 f''(\theta) \Big|_{\text{stationary}} &= -k \left[\left(\frac{k'}{k}\right)' - 1 - \frac{k'^2}{k^2} \right] \cos(\theta - \varphi) \\
 &= \pm k \left[\left(\frac{k'}{k}\right)' - 1 - \frac{k'^2}{k^2} \right] / \sqrt{1 + \left(\frac{k'}{k}\right)^2} \\
 &= \pm k \left(\frac{\omega_j + \omega_{j+1}'}{\omega_j - \omega_{j+1}} \right) \frac{1 - \frac{1}{2} \left(\frac{k'}{k}\right)^2}{\sqrt{1 + \left(\frac{k'}{k}\right)^2}} \quad , \quad (A.15)
 \end{aligned}$$

$$2.12 \quad \frac{\omega_1 + \omega_2}{\omega_2 - \omega_1} = \frac{1}{\sqrt{1 - \frac{4\omega_1 \omega_2}{g} \cos \theta}} \quad ,$$

$$\begin{aligned}
 f'''(\theta) &= (k''' - k') \cos(\theta - \varphi) - (k'' - k) \sin(\theta - \varphi) - 2k'' \sin(\theta - \varphi) \\
 &\quad - 2k' \cos(\theta - \varphi) \\
 &= (k''' - 3k') \cos(\theta - \varphi) - (3k'' - k) \sin(\theta - \varphi), \quad (A-16)
 \end{aligned}$$

$$\begin{aligned}
 f'''(\theta) \Big|_{\theta=0} &= [k''' - 3k' - \frac{(3k'' - k)k'}{k}] \cos(\theta - \varphi) = \frac{k k''' - 3k' k'' + k k' - 3k k''}{k} \cos(\theta - \varphi) \\
 &= (k''' - 2k' - 3 \frac{k' k''}{k}) \cos(\theta - \varphi) = \left[k \left(\frac{k'}{k} \right)'' - 2k' - \frac{2k' k''}{k} \right] \cos(\theta - \varphi) \\
 &= k \left[\left(\frac{k'}{k} \right)'' - 2 \frac{k'}{k} \left(1 + \frac{k''}{k} \right) \right] \cos(\theta - \varphi), \quad (A-17) \\
 &= k \left[\left\{ \left(\frac{k'}{k} \right)' - \left(1 + \frac{k''}{k} \right) \right\}' \right] \cos(\theta - \varphi)
 \end{aligned}$$